- Ask questions: let me know how you perceive nature so that if there is a misperception, we can correct it
- Ask simple questions
- if you have any doubt, repeat what you have understood. It need not be in the form of a question
- Don't try to guess the type of questions that I can ask, try to understand nature
- Bad habits that you have learned in years takes more than a few months to correct!

Department of Physics

The only forces that we will study in this year are gravity, and EM force. EM force represents itself through

- Friction
- Any pull or a push (e.g. Normal Force, forced due to the tension on a string, force acting by a spring)



TRACKER Software (http://www.cabrillo.edu/ dbrown/tracker/)



Altuğ Özpineci (METU)

Phys109-MECHANICS

Review of Work Done By Gravity

- $W_{tot} = \Delta T$ where $T = \frac{1}{2}mv^2$
- *W_G* = −*mg*∆*h* depends only on the height difference, and on nothing else
- If an object moves under the influence of gravity only then throughout the motion

$$\frac{1}{2}mv^2 + mgh = const \tag{97}$$



Work Done By Gravity

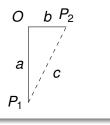
- The force acting on an object of mass *m* is $\vec{F}_w = -mg\hat{z}$ (\hat{z} points upwards)
- If the object is displaced by $d\vec{\ell}$, then $dW = \vec{F}_w \cdot d\vec{\ell} = -mgdz$, i.e. the work done by its weight is proportional to the change in its *z* coordinate (its height)
- As the object goes from P_1 to P_2 , the to calculate the total work, just sum the changes in its height. Hence total work is $W_{tot} = -mg\Delta h$

Department of Physics

PHYS109

128 / 141

Work Done By Friction



Altuğ Özpineci (METU)

- Assume a constant friction force of magnitude *F*_f.
- *a*, *b* and *c* are the corresponding length.

•
$$W_{P_1 \rightarrow O} = -F_f a$$

•
$$W_{O \rightarrow P_2} = -F_f b$$

•
$$W_{P_1 \rightarrow O \rightarrow P_2} = -F_f(a+b)$$

•
$$W_{P_1 \rightarrow P_2} = -F_f c \neq -F_f(a+b)$$

 Hence the work done by friction depends on how one goes from the initial point to the final point

 Department of physics

 Fizik

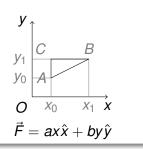
 Böllimü

 Böllimü

 Böllimü

 Phys109-MECHANICS
 PHYS109
 129 / 141

Example: $\vec{F} = ax\hat{x} + by\hat{y}$



Work done by the force as one goes from *A* to *B* through *C*:

$$W_{A \to C \to B} = W_{A \to C} + W_{C \to B}$$

• Along the path $A \to C$, $x = x_0$ and
hence $\vec{F} = ax_0\hat{x} + by\hat{y}$, $d\vec{\ell} = dy\hat{y}$
• $dW = \vec{F} \cdot d\vec{\ell} = bydy$

$$\int_0^{W_{A\to C}} dW = \int_{y_0}^{y_1} by dy \qquad (98)$$

$$W_{A\to C} = \frac{1}{2}by_1^2 - \frac{1}{2}by_0^2 \quad (99)$$



Altuğ Özpineci (METU)

Example: $\vec{F} = ax\hat{x} + by\hat{y}$

 Work done by the force as one goes from A to B through C: $W_{A \to C \to B} = W_{A \to C} + W_{C \to B}$ • Along the path $A \rightarrow C$, $y = y_1$ and hence $\vec{F} = ax\hat{x} + by_1\hat{y}, d\vec{\ell} = dx\hat{x}$ • $dW = \vec{F} \cdot d\vec{\ell} = axdx$ $\int_{a}^{W_{C\to B}} dW = \int_{x_{a}}^{x_{1}} ax dx$ (98) $W_{C \to B} = \frac{1}{2}ax_1^2 - \frac{1}{2}ax_0^2$ (99)



Altuğ Özpineci (METU)

Phys109-MECHANICS

Work Done by Friction

Example: $\vec{F} = ax\hat{x} + by\hat{y}$

 y_{1} y_{1} y_{0} A A X_{1} X $\vec{F} = ax\hat{x} + by\hat{y}$ $W_{A \to C} = \frac{1}{2}by_{1}^{2} - \frac{1}{2}by_{0}^{2}$ $W_{C \to B} = \frac{1}{2}ax_{1}^{2} - \frac{1}{2}ax_{0}^{2}$

Work done by the force as one goes from
A to *B* through *C*:

$$W_{A \to C \to B} = W_{A \to C} + W_{C \to B}$$

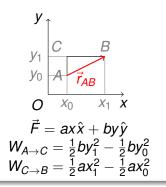
 $W_{A \to C \to B} =$
 $(\frac{1}{2}ax_1^2 + \frac{1}{2}by_1^2) - (\frac{1}{2}ax_0^2 + \frac{1}{2}by_0^2)$

Department Department Bölümü Physics Bölümü C D C C D C C D C C D C C D C C D C C D C C D C C D C C D C C D C

Altuğ Özpineci (METU)

Phys109-MECHANICS

Example: $\vec{F} = ax\hat{x} + by\hat{y}$



Work done by the force as one goes from *A* to *B* through a straight line:

- Let $\vec{r}_A = x_0 \hat{x} + y_0 \hat{y}$, $\vec{r}_B = x_1 \hat{x} + y_1 \hat{y}$, $\vec{r}_{AB} = \vec{r}_B - \vec{r}_A$
- Any point on the trajectory can be written as r
 [−] = r
 [−] + λr
 [−]
- $d\vec{\ell} = (d\lambda)\vec{r}_{AB},$ $dW = \vec{F} \cdot d\vec{\ell} = d\lambda\vec{F} \cdot \vec{r}_{AB} =$ $d\lambda(a(x_1\lambda + x_0(1 - \lambda))(x_1 - x_0) +$ $b(y_1\lambda + y_0(1 - \lambda)(y_1 - y_0))$
- $W_{A \to B} = \int_0^{W_{A \to B}} dW = \int_0^1 d\lambda(\cdots) = (\frac{1}{2}ax_1^2 + \frac{1}{2}by_1^2) (\frac{1}{2}ax_0^2 + \frac{1}{2}by_0^2)$

< ロ > < 同 > < 回 > < 回 >



Altuğ Özpineci (METU)

PHYS109 130 / 141

Example: $\vec{F} = ax\hat{x} + by\hat{y}$

 y_{1} y_{1} y_{0} $F = ax\hat{x} + by\hat{y}$ $W_{A \to C} = \frac{1}{2}by_{1}^{2} - \frac{1}{2}by_{0}^{2}$ $W_{C \to B} = \frac{1}{2}ax_{1}^{2} - \frac{1}{2}ax_{0}^{2}$

Work done by the force as one goes from *A* to *B* through a straight line:

- Let $\vec{r}_A = x_0 \hat{x} + y_0 \hat{y}$, $\vec{r}_B = x_1 \hat{x} + y_1 \hat{y}$, $\vec{r}_{AB} = \vec{r}_B - \vec{r}_A$
- Any point on the trajectory can be written as r
 [−] = r
 [−] + λr
 [−] μ
 [−] (λ ≤ 1)

•
$$d\vec{\ell} = (d\lambda)\vec{r}_{AB},$$

 $dW = \vec{F} \cdot d\vec{\ell} = d\lambda\vec{F} \cdot \vec{r}_{AB} =$
 $d\lambda(a(x_1\lambda + x_0(1 - \lambda))(x_1 - x_0) +$
 $b(y_1\lambda + y_0(1 - \lambda)(y_1 - y_0)))$
• $W_{A \rightarrow B} = \int_0^{W_{A \rightarrow B}} dW = \int_0^1 d\lambda(\cdots) =$
 $(\frac{1}{2}ax_1^2 + \frac{1}{2}by_1^2) - (\frac{1}{2}ax_0^2 + \frac{1}{2}by_0^2)$
from A to B is the same in both paths. And
 $M(x - y_0) = M(x - y_0)$

Work done as one goes from *A* to *B* is the same in both paths. And the can be written as $W_{A \rightarrow B} = U(x_0, y_0) - U(x_1, y_1)$

Conservative Forces

- **Definition:** A force is conservative if the work done by that force as an object moves from a point *P*₁ to a point *P*₂ is independent of the path that the object takes.
- Friction is an example of a *non-conservative* force.
- Gravity (any constant force in general) is a *conservative* force.
- $\vec{F} = ax\hat{x} + by\hat{y}$ is a conservative force

Departmen of Physics

• • • • • • • • • • • •

- If \vec{F} is conservative $W = \int_{P_1}^{P_2} \vec{F} \cdot d\vec{\ell}$ is independent of how one goes from P_1 to P_2 .
- To evaluate *W*, one can *choose* any path
- Define a function U(P) such that $U(P) = U(P_0) \int_{P_0}^{P_1} \vec{F} \cdot d\vec{\ell}$.
- $W = \Delta T$
- Suppose the object moves from P_1 to P_2 under the influence of the conservative force \vec{F} .
- The work done by \vec{F} is:

$$W = \int_{P_1}^{P_2} \vec{F} \cdot d\vec{\ell} = \int_{P_1}^{P_0} \vec{F} \cdot d\vec{\ell} + \int_{P_0}^{P_2} \vec{F} \cdot d\vec{\ell}$$
(98)
= $(U(P_1) - U(P_0)) + (U(P_0) - U(P_2)) = U(P_1) - U(P_2)$ (99)
• $W = T_2 - T_1$:

$$T_2 - T_1 = U(P_1) - U(P_2)$$

$$T_1 + U(P_1) = T_2 + U(P_2)$$
Department of Physics

Conservation of Energy

- U is called the potential energy
- T + U is called the mechanical energy
- For an object moving under the influence of a conservative force only T + U is always conserved.
- An object that is raised by h, has *potential* to do work, it has a larger *potential* energy
- Potential energy is NOT a property of a single object, but a property of the system as a whole.

Departmen of Physics

• • • • • • • • • • • •

- Newton's three laws are vector relations
- Conservation of energy is a scalar relation
- If the expression for potential energy is known, speed at any point can be determined without solving any differential equation or integrals.



Force from Potential

- Consider two nearby points separated by the vector $d\vec{r}$.
- The difference in their potential energies is:

$$U(\vec{r}+d\vec{r})-U(\vec{r})=-\vec{F}(\vec{r})\cdot d\vec{r}$$
(102)

• Above is valid for any
$$d\vec{r}$$
. Denote $U(\vec{r}) \equiv U(x, y, z)$ if $\vec{r} = x\hat{x} + y\hat{y} + z\hat{z}$

• If $d\vec{r} = dx\hat{x}$, then

$$U(x + dx, y, z) - U(x, y, z) = -F_x(x, y, z)dx$$
 (103)

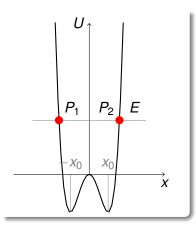
$$\implies F_x(x,y,z) = -\frac{U(x+dx,y,z) - U(x,y,z)}{(x+dx) - x} \equiv -\frac{\partial U}{\partial x} \quad (104)$$

• Similarly $F_y = -\frac{\partial U}{\partial y}$ and $F_z = -\frac{\partial U}{\partial z}$ • Hence $\vec{F} = -(\hat{x}\frac{\partial U}{\partial x} + \hat{y}\frac{\partial U}{\partial y} + \hat{z}\frac{\partial U}{\partial z}) \equiv -\vec{\nabla}U$ where $\vec{\nabla}$ is the *nabla* operator.

Altuğ Özpineci (METU)

PHYS109 135 / 141

Interpreting Potential Graphs

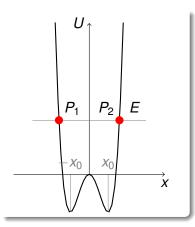


- Consider one dimensional example $F_x = -\frac{dU}{dx}$, i.e. F_x points in the direction that U(x) is decreasing
- Consider a particle with total energy E. Then K = E - U > 0, i.e. The particle can only be at points for which U(x) < E.
- points x such that U(x) = E are called turning points



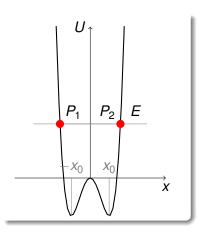
Altuă Özpineci (METU) 136 / 141

Interpreting Potential Graphs



- At points $x = \pm x_0$, and x = 0, the force acting on an object is zero: they are called equilibrium points.
- If an object is slightly displaced from x = ±x₀, they try to move towards ±x₀: they are called stable equilibrium points
- If an object is slightly displaced from x = 0, they try to move away from x = 0: x = 0 is an unstable equilibrium point.

Interpreting Potential Graphs



- If an object is slightly displaced from x = ±x₀, they try to move towards ±x₀: they are called stable equilibrium points
- If an object is slightly displaced from x = 0, they try to move away from x = 0: x = 0 is an unstable equilibrium point.
- If an object is displaced slightly from an equilibrium point, it neither goes towards nor away from the equilibrium point, it is called neutral equilibrium point.

Dissipative Forces

- Dissipative forces in fact convert mechanical energy to internal energy
- The total energy of the universe is constant
- $\Delta(T + U) = W_{dissipative forces}$



Gravitational Potential Energy and Escape Velocity

• Gravitational force acting on a mass *m*, due to a mass *M* is:

$$\vec{F} = -G_N \frac{mM}{r^2} \hat{r} \tag{105}$$

• The work done when *m* is displaced by $d\vec{\ell}$ is

$$dW = \vec{F} \cdot d\vec{r} = -G_N \frac{mM}{r^2} \hat{r} \cdot d\vec{\ell} = -G_N \frac{mM}{r^2} dr \qquad (106)$$

where dr is the change in the radial distance, i.e. radial component if $d\vec{\ell}$.

Potential energy difference is

$$U(P) - U(\infty) = -\int_{\infty}^{P} \left(-G_{N} \frac{mM}{r^{2}} dr\right)$$
(107)
$$= -G_{N} \left.\frac{mM}{r}\right|_{r=\infty}^{r=r_{P}} = -G_{N} \frac{mM}{r_{P}}$$

• In general $U(\infty)$ is chosen to be zero $U(\infty) = 0$
• In general $U(\infty)$ is chosen to be zero $U(\infty) = 0$
• Hysios PHYSIOS (METU) Physios PHYSIOS (METU) (188/141)

Gravitational Potential Energy and Escape Velocity

Q: What should be the initial speed of an object on the surface of a planet of mass *M* and radius *R*, if it is to go until infinity?

- Let v_0 and v_∞ be the initial and final speeds.
- Initial mechanical energy is $E = \frac{1}{2}mv_0^2 + G_N \frac{mM}{R}$.
- Final mechanical energy is $E = \frac{1}{2}mv_{\infty}^2$
- Conservation of mechanical energy:

$$\frac{1}{2}mv_0^2 - G_N \frac{mM}{R} = \frac{1}{2}mv_\infty^2$$
(109)
$$\Rightarrow v_0^2 = \frac{2}{m} \left(\frac{1}{2}mv_\infty^2 + G_N \frac{mM}{R}\right)$$
(110)

• The minimum possible speed is called the escape velocity:

$$v_{esc} = \sqrt{G_N \frac{2M}{R}}; v_{esc,Earth} = 11.2 \text{ km/s} = 40320 \text{ km}/\text{physics}$$

Potential Energy of a Spring

- The work done by a spring on an object as it moves from x = 0 to x = L was calculated as $W = -\frac{1}{2}kL^2$.
- The potential energy of a spring that is stretched by *L* is $U(L) = \frac{1}{2}kL^2$

